

Appendix M

Reciprocity for the Bidirectional Reflectance

In this appendix we prove the *Principle of Reciprocity* for the bidirectional reflectance, that is

$$\rho(\nu; \theta', \phi'; \theta, \phi) = \rho(\nu; \theta, \phi; \theta', \phi'). \quad (\text{M.1})$$

Referring to Fig. M.1, the proof first determines the exchange of radiative energy between the black elements dA_1 and dA_2 in a *hohlraum* due to reflection of energy by the surface dA_2 . The theorem is proven by equating the energy exchange $dE_{\nu 123}$ from 1 to 3 via 2, and the energy

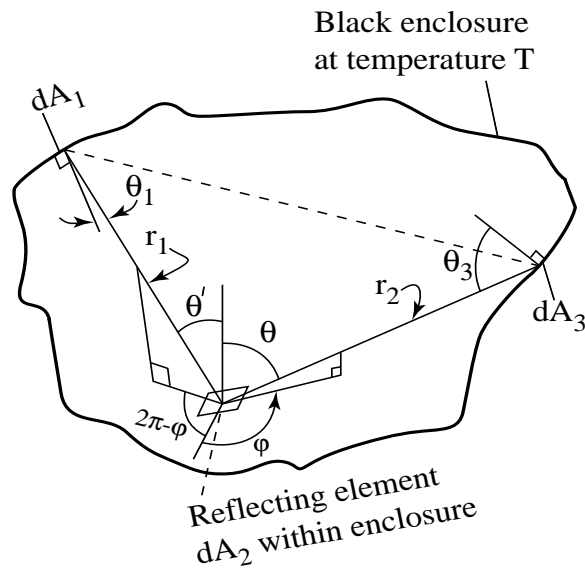


Fig. AM.1. Exchange of radiative energy within a *hohlraum*.

exchange $dE_{\nu 321}$ from 3 to 1 via 2, and equating these two quantities (see Figure AM.1).

The energy exchange $dE_{\nu 123}$ must balance the reciprocal energy exchange $dE_{\nu 321}$ in a TE situation. Otherwise there would be a net heating/cooling of one of the areas at the expense of the other, and this violates the conditions of the *hohlraum*. Let's first consider the exchange from 1 to 3 via 2. The radiative energy reflected by dA_2 and intercepted by dA_3 is given by that 'emitted' into the solid angle subtended by dA_3 , $dA_3 \cos \theta_3 / r_2^2$,

$$dE_{\nu 123} = I_{\nu r_2}^+(\theta', \phi'; \theta, \phi) \cos \theta \frac{dA_3 \cos \theta_3}{r_2^2} d\nu dt. \quad (\text{M.2})$$

The reflected intensity $I_{\nu r_2}^+$ is related to the intensity $I_{\nu}^-(\theta', \phi')$ arriving at dA_2 through eqn. 5.15

$$I_{\nu r_2}^-(\theta', \phi') = \cos \theta' \rho(\nu; \theta', \phi'; \theta, \phi) I_{\nu 1}^-(\theta', \phi') d\omega_{21} \quad (\text{M.3})$$

where

$$d\omega_{21} = \frac{dA_1}{r_1^2} \cos \theta_1 \quad (\text{M.4})$$

is the angle subtended by dA_1 from the point dA_2 . Putting these together, we find that the rate of energy exchange per unit frequency from 1 to 3 via 2 is

$$\frac{dE_{\nu 123}}{d\nu dt} = I_{\nu 1}^-(\theta', \phi') \rho(\nu; \theta', \phi'; \theta, \phi) \cos \theta' \left[\frac{dA_1}{r_1^2} \cos \theta_1 \right] \cos \theta \left[\frac{dA_3}{r_2^2} \cos \theta_3 \right]. \quad (\text{M.5})$$

Now consider energy exchange in the reverse direction, 3 to 1 via 2. The rate of energy per unit frequency reflected at dA_2 into the direction of dA_1 is

$$\frac{dE_{\nu 321}}{d\nu dt} = I_{\nu r_2}^+(\theta, \phi; \theta', \phi') \cos \theta' dA_2 \left[\frac{dA_1}{r_1^2} \cos \theta_1 \right]. \quad (\text{M.6})$$

But the reflected intensity at dA_2 is given by

$$I_{\nu r_2}^+(\theta, \phi; \theta', \phi') = I_{\nu 3}^-(\theta, \phi) \rho(\nu; \theta, \phi; \theta', \phi') \cos \theta d\omega_{23} \quad (\text{M.7})$$

where $d\omega_{23} = dA_3 \cos \theta_3 / r_2^2$ is the solid angle subtended by dA_3 at the point dA_2 . Putting these together, we find for the energy exchange rate

$$\frac{dE_{\nu 321}}{d\nu dt} = I_{\nu 3}^-(\theta, \phi) \rho(\nu; \theta, \phi; \theta', \phi') \cos \theta' \cos \theta \left[\frac{dA_3}{r_2^2} \cos \theta_3 \right] \left[\frac{dA_1}{r_1^2} \cos \theta_1 \right]. \quad (\text{M.8})$$

Equating the two rates of energy exchange, we find

$$I_{\nu 1}^-(\theta', \phi') \rho(\nu; \theta', \phi'; \theta, \phi) = I_{\nu 3}^-(\theta, \phi) \rho(\nu; \theta, \phi; \theta', \phi'). \quad (\text{M.9})$$

But in TE, the two intensities are just the Planck function, $I_{\nu 1}^- = I_{\nu 3}^- = B_\nu$. Therefore

$$\rho(\nu; \theta', \phi'; \theta, \phi) = \rho(\nu; \theta, \phi; \theta', \phi'). \quad (\text{M.10})$$

Given the above reciprocity property for the BDRF, we now show that reciprocity also applies to the flux reflectance. Placing the two definitions together, we have

$$\rho(\nu; -\hat{\Omega}', 2\pi) = \int_+ d\omega \cos \theta \rho(\nu; -\hat{\Omega}', \hat{\Omega}) \quad (\text{M.11})$$

$$\rho(\nu; 2\pi, \hat{\Omega}) = \int_- d\omega' \cos \theta' \rho(\nu; -\hat{\Omega}', \hat{\Omega}). \quad (\text{M.12})$$

The first quantity, $\rho(\nu; -\hat{\Omega}', 2\pi)$ is the *directional-hemispherical reflectance*, and the second quantity, $\rho(\nu; 2\pi, \hat{\Omega})$ is the *hemispherical-directional reflectance*. If we evaluate the first of the above equations at $\hat{\Omega}' = \hat{\Omega}$, and place primes on the angular integration variables (realizing that they are dummy variables), we have

$$\rho(\nu; -\hat{\Omega}, 2\pi) = \int_+ d\omega' \cos \theta' \rho(\nu; -\hat{\Omega}, \hat{\Omega}'). \quad (\text{M.13})$$

Invoking reciprocity of the BDRF, $\rho(\nu; -\hat{\Omega}, \hat{\Omega}') = \rho(\nu; -\hat{\Omega}', \hat{\Omega})$, we have

$$\rho(\nu; -\hat{\Omega}, 2\pi) = \int_+ d\omega' \cos \theta' \rho(\nu; -\hat{\Omega}', \hat{\Omega}). \quad (\text{M.14})$$

But this is the same expression for the hemispherical-directional reflectance, eqn. M.12. Thus we find the desired reciprocity relationship

$$\rho(\nu; -\hat{\Omega}, 2\pi) = \rho(\nu; 2\pi; \hat{\Omega}). \quad (\text{M.15})$$